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The meaning of the speed of sound at which sound waves travel through a medium (usually air) was derived from the Latin concept "celerita", meaning "velocity"⁹.

Historically speaking the speed of sound was one of, if not the first effects relating to the study of acoustics and variables relating to sound propagation outdoors. Due to the fact that the speed of sound is one of the most historically significant factors in acoustics and sound propagation this section of the report discusses a brief history essential to the subject.

It is possible that as long ago as the 6th century BC, the Greek mathematician Pythagoras was aware that sound was a vibration transmitted from a source to the ear.

Written records dating from the time of the Greek philosopher Aristotle (350 BC)¹⁰ show a rudimentary knowledge of sound propagation, as he wrote:

"All sounds are produced because air is set in motion by expansion or compression or when it is clashed together by an impact from the breath or from the strings of musical instruments"

Pythagoras thought that high frequencies were transmitted through air more rapidly than low frequencies. In addition to Pythagoras comments, the Roman architect Vitruvius (25 BC)¹¹ wrote:

"Voice is breath, flowing and made sensible to the hearing by striking the air. It moves in infinite circumferences of circles as when, by throwing a stone into

still water, you produce innumerable circles of waves increasing from the centre and spreading outwards."

The first explicit mention of sound velocity appears to be by Francis Bacon¹². He discussed the possibility of comparing the velocity of sound with that of light (which he knew at the time to be immeasurably high) by comparing the time taken for the sound of a church bell to travel one mile (1.6km) with that taken by a simultaneous light signal (an interrupted lighted taper) over the same distance by using his own pulse as the timing mechanism.

However, the first actual quantitative determination was by the French mathematician Mersenne¹³, working under the influence of Galileo.

Mersenne used both musical instruments and gun fire as sound sources and estimated the distance travelled by sound in one second to be equal to 230 French Toises; this corresponds to 448 m/s and is thus higher than the known true value.

Mersenne incorrectly asserted that the same speed was observed by night and by day, either with wind or against it¹⁴.

By comparing the speed of travel of sound produced by a large weapon such as a cannon with that produced by a small weapon such as a musket, Gassendi¹⁵ demonstrated the "surprising" fact that the time taken for sound to traverse a given distance is independent of both pitch and intensity, but as with Mersenne he erroneously concluded

that wind had no effect on the velocity of sound. (Refer to Lenihan¹⁶ for a discussion of the roles of Mersenne and Gassendi in determining the velocity of sound.)

The importance of air as the medium through which sound is transmitted had yet to be established. The far-reaching work on barometric pressure by Torricelli¹⁷ led to his successful demonstration of a vacuum and soon a number of experiments with air pumps were under way¹⁸. During the 17th century Sir Isaac Newton¹⁹ showed that in an elastic fluid the speed at which sound waves are propagated is proportional to the square root of the elasticity divided by density. In fact, the reasoning which led to this result was so obscure that few have claimed to be able to follow it.

The noted mathematicians d'Alembert and Bernoulli both concluded that this was the most obscure and difficult part of the whole of Newton's Principia, whilst at one time Lagrange²⁰ actually claimed the derivation to be illogical.

Illogical or obscure it may have been, the work marked an important step forward and Newton went on to use Boyle's Law, which holds true only for constant temperature, to derive the speed of sound in air as $c = \sqrt{p/\rho}$, leading to a calculated value of 968 ft/s, or approximately 295 m/s.

Newton concluded that experimental data then available indicated that the velocity lay between 280 and 330 m/s, and in

the first edition of Principia he appears satisfied with the order of magnitude agreement between theory and experiment.

A number of experimental determinations of the speed of sound were reported about this time but a particularly detailed study including a review of previous determinations was reported by the Reverend William Derham in 1708. He arranged for guns (called sakers) to be fired from various church towers and other eminences in the neighbourhood, covering distances of up to 12.5 miles (20.1 km) from his own church tower at Upminster, observed the flash (in daylight using a telescope for this purpose) and timed the interval to the arrival of the report using an accurate portable movement with pendulum beating half seconds. In this way he confirmed that velocity was independent of level (i.e. independent of distance from the source) and arrived at an accurate mean value of 348 m/s²¹.

In a further series of tests over a fixed distance of 12.5 miles Derham investigated various effects on the velocity of sound and, contrary to suggestions by previous workers, Derham concluded that favourable winds accelerated sound propagation whilst opposing winds retarded it.

Judged by today's standards his results appear highly creditable and remarkable accurate. Unfortunately it is not certain that he established the correct quantitative relation between effective sound speed and wind speed, for this was before the invention of the anemometer (although the ingenious experimenter Hooke had in fact invented a swinging-plate wind velocity indicator in 1667).

Derham's estimates of wind speed were largely based on an arbitrary fifteen point scale and attempts to translate to true wind speed had not been successful²².

Although graduated air thermometers had been in use since the time of Galileo, thermometry was still being established; Fahrenheit's scale based on mercury in glass thermometer was not published until 1724²³ although probably established around 1714, the Celsius scale not until 1724²⁴.

Derham did not actually measure the prevailing temperature but concluded that the speed of sound was the same in winter as in summer and, although this false conclusion has been held against him, this was surely but a major lapse.

In 1738 Cassini and others from the Académie des Sciences²⁵ made measurements over a distance of

approximately 28km. They used cannons as sources, pendulum clocks for timing, and worked at night so that stable meteorological conditions prevailed²⁶. They were the first to state definitively that with a prevailing wind speed denoted as "u", in the direction of the wind sound is propagated with speed (c + u) but that against the wind the effective speed of sound is (c - u).

They used reciprocal firing from either end of the baseline and adopted the mean traverse time in order to minimise intrusive wind effects and arrived at a value of 337 m/s for the speed of sound²⁷.

Although Cassini concluded that temperature does affect the velocity of sound, the first quantitative study of the influence of temperature was undertaken by Bianconi²⁸ at Bologna by comparing velocities determined in winter and in summer. It was correctly concluded that an increase in temperature produces an increase in the velocity of sound.

The remainder of the 18th Century was essentially a period of consolidation. In view of the difficulties encountered with Newton's derivation, several theoreticians tackled the problem and a reasonably clear treatment was eventually produced by Euler²⁹, whilst Lagrange revised Newton's

sound weighted standardized impact sound pressure levels structure born sound low frequency noise octave band time weighting sabin speech intelligibility noise reduction engineering sound level environment spectrum resource management SIL ambient sound insulation vibration rumble sound level meter noise map silencer emission speaker amenity value

reverberation time noise reduction coefficient Dntw speech transmission index dBA frequency band noise Hertz or Hz far field octave airborne sound impact sound pressure level immission plane wave SEL line source random incidence sound reduction index.

R best practical option frequency spectrum noise exchange rate logarithm live room limiter calibration room criterion curves habitat structure sound power sound

pressure level hiss free field Ctr articulation class ambience Bel acoustics environment assessment structural analysis apparent sound reduction index resonance natural frequency flow kinetic measurement prediction signal processing threshold shift shadow zone transducer wavelength narrow band overtone reflection percentile level impedance directivity fresnel number harmonic echo ambient active noise control attenuation coverage angle coincidence hearing point abatement temperature diffusion indoors reflections concave node anti-node wind

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reasoning and generalised the treatment to cover sound waves of arbitrary character (i.e. not just simple harmonic waves). However, all the calculated values were seriously discrepant from the experimental data³⁰.

It was not until Laplace³¹, following the observation by Dalton that a sudden compression produces heating of a gas, pointed out that setting the volume elasticity of the air equal to the static pressure implied that isothermal conditions existed, whilst due to the rapidity of the pressure fluctuations, adiabatic (an adiabatic process is a process in which no heat is gained or lost in the working fluid) conditions probably prevailed.

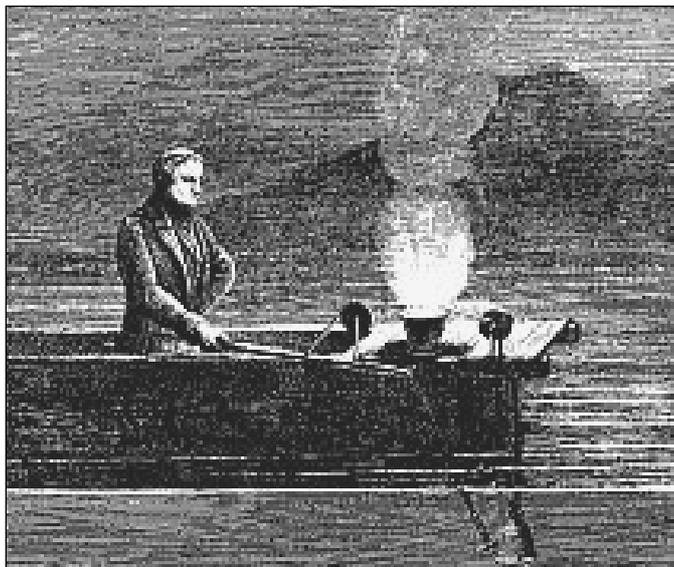
This led to the revised relation $c = \sqrt{\gamma p / \rho}$, where γ is the ratio of the specific heat at constant pressure to that at constant volume.

It is understood the inclusion of γ immediately brought calculated and observed values for the velocity of sound into substantial agreement. Although at this stage it was not entirely obvious that sound propagation took place completely adiabatically.

Any lingering doubts were largely resolved when Stokes³² showed that

sound propagation must occur either substantially adiabatically or substantially isothermally, as otherwise large damping factors would prevail and a comparatively high attenuation rate would be observed³³.

As Herzfeld and Rice³⁴ have shown, in an unbounded wave the time required for temperature



Measurement of the propagation speed of sound in water of the Lake Geneva in 1828 (observation point 1) according Guillemain (1868)

equilibration to occur is proportional to the square of the wave length, whereas the time actually available for this to occur is only directly proportional to wave length. Thus propagation becomes more truly adiabatic as the sound frequency decreases whilst, at normal atmospheric pressure,

adiabaticity is expected to persist up to the highest attainable frequencies³⁵.

A definitive determination of sound velocity was undertaken in 1822 by a commission appointed by the Bureau des Longitudes which included Prony, Arago, Bouvard, Mathieu, Gay-Lussac and Humboldt³⁶. They used accurate chronometers, an accurately surveyed distance of 18.6223km and again used reciprocal firing of cannons to minimise wind effects. They arrived at a mean time for sound to travel the measured distance of 54.63 s at 16 °C, leading to a speed value of approximately $c_0 = 331.2$ m/s.

As time passed, other techniques for determining the velocity of sound were put forward. It had now become obvious that in the open air the accuracy of speed determination, particularly when made over the long base lengths necessary to ensure adequate time resolution, was limited by uncertainties regarding the temperature and the wind velocity.

For instance a relatively low wind speed of 5 miles per hour (8km/hr) corresponds to a possible error of 2.2 m/s; although a substantial part



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of this error could be eliminated by reciprocal firing, at least part would remain unless exact synchrony were achieved.

Also, a change of temperature of 1°C produces a speed change of about 0.6 m/s. Besides carrying out open-air measurements at Versailles, Regnault³⁷ therefore carried out an elaborate series of experiments in the newly-laid water pipes under the city of Paris using sources consisting of pistols, explosions, and musical instruments. Now when confined in pipes the speed of sound is less than in free space but by using pipes of different diameter he managed to extrapolate to free-field conditions and arrived at a value for dry air of $c_0 = 330.7 \text{ m/s}$ ³⁸.

Tackling the problem in an entirely different way Hebb³⁹, while working at the Ryerson Physical Laboratory in the USA, made direct measurement in free air of the wavelength of sound corresponding to a signal of known frequency. The technique, which had been proposed by Michelson, is said to be almost too well known to require description.

Two paraboloidal reflectors (focal length approximately 0.38 m, diameter 1.5 m) were arranged coaxially, one of them being moveable on a parallel track. A pure-tone air-whistle source whose frequency $f = 2376.5 \text{ Hz}$ was found by comparison with a tuning fork, was placed at the focus of one of the mirrors whilst a carbon microphone was placed at the focus of each of the mirrors. Using a split-primary transformer the outputs of the two microphones were combined in such a way that when the output from the secondary was monitored using headphones the sound heard was proportional to the vector sum of the two microphone outputs.

By progressively changing the separation between the mirrors the relative phases of the two signals varied so that at certain points they cancelled whilst at other points they reinforced each other. This method

therefore provided a direct determination of wavelength and, as the frequency was known, this yielded the sound velocity.

The experiment was carried out in a hall 36m long, thus no wind was encountered and the temperature was found to remain very constant. The separation between mirrors could be increased to as much as 100λ and the positions of the minima could be located to within about 1 cm, so that an accuracy of the order of 0.1% was achieved. In this way Hebb reported a value of $c_0 = (331.29 \pm 0.04) \text{ m/s}$ but subsequently⁴⁰ he revealed that this result was slightly in error due to the method he had used to correct the velocity measured in air containing moisture to obtain the dry-air value.

His revised estimate for dry air was remarkable $c_0 = 331.41 \text{ m/s}$.

The development of instrumentation during World War 1 for use in connection with sound ranging of enemy guns, notably the hot-wire microphone in the neck of Helmholtz Resonator developed by Tucker, led to further open-air determinations of sound velocity.

Among these were experiments reported by Esclançon⁴² and a particularly detailed study by Angerer and Ladenburg⁴³, but the day of open-air measurements was long since past and their

determinations were doubtless more accurate as methods for measuring the mean temperature along the sound path than as primary determinations of speed of sound⁴⁴.

Measurements of velocity of sound using a piezoelectric transducer and interferometric technique were reported by Pierce⁴⁵ who appears to be the first to report velocity dispersion (i.e. frequency-dependent velocity of sound) in air, based primarily on the laboratory studies and reviews of previous work reported by Hardy, Telfair and Pielemeier⁴⁶ supplemented by the work of Smith⁴⁷.

At 1 kHz and 1 atmospheric pressure the velocity of unbounded progressive plane waves in dry air containing 0.03% carbon dioxide and at 0 °C is given as $(331.4 \pm 0.05) \text{ m/s}$.

Harris⁴⁸, who made careful measurements of relative sound velocity, has shown that at room temperature it reaches a minimum value at a relative humidity of about 14% at which point it is approximately 0.2 m/s lower than the dry-air value: at 100% Relative Humidity (RH) the velocity is about 1.1 m/s higher than the dry-air value⁴⁹. It is this value of 331m/s which is commonly referenced to in text books today.

The speed (velocity =speed with

| Substance | Temp (°C) | Speed (m/sec) | Speed (ft/sec) |
|-----------------|-----------|---------------|----------------|
| CO ₂ | 0 | 258 | 816 |
| CO ₂ | 35 | 274 | 900 |
| Air | 0 | 331.5 | 1,087 |
| Air | 20 | 344 | 1,130 |
| Water Vapor | 35 | 402 | 1,320 |
| Helium | 20 | 927 | 3,040 |
| Hydrogen | 0 | 1,270 | 4,165 |
| Water | 15 | 1,437 | 4,714 |
| Steel | 20 | 5,000 | 16,400 |

Figure 22: Speed of sound in various substances and temperatures.

direction) of sound in dry air is generally given as:

$$V_{\text{sound in air}} \approx 331.6 + 0.6 T_c \text{ (m/s)}$$

(Equation 20)

The speed of waves is affected by the density of the medium and the density is affected by the temperature and pressure, the speed of sound through air varies depending on these (and other) factors.

Bieler-Butticaz⁵⁰ appears to have been among the first to note that the attenuation of sound in air is strongly dependent upon temperature and humidity although she gave no quantitative data. The speed varies depending on atmospheric conditions; the most important factor is the temperature.

Bieler-Butticaz noted:

“The humidity has very little effect on the speed of sound, while the static sound pressure (air pressure) has none. Sound travels slower with an increased altitude (elevation if you are on solid earth), primarily as a result of temperature and humidity changes.”

Einstein⁵¹ discussed the velocity of sound in mixtures of dissociated and non-dissociated molecules of a diatomic gas and showed that it is possible to calculate the rate of energy transfer between the two kinds of molecule from a determination of the velocity of sound as a function of frequency.

This concept was applied to the calculation of sound attenuation in a gas by Herzfeld and Rice and soon afterwards Kneser⁵² showed their analysis adequately explained the sound dispersion previously observed in carbon dioxide.

In a non-dispersive medium the speed of sound is independent of frequency, therefore the speed of energy transport and sound propagation are the same. Air is a non-dispersive medium. In a dispersive medium the speed of sound is a function of frequency.

The spatial and temporal distribution of a propagating disturbance will continually change. Each frequency component propagates at each its own phase speed, while the energy of the disturbance propagates at the group velocity.

Water is an example of a dispersive medium.

The speed of sound (c) in a medium depends on the medium's elasticity (E) and its density (ρ) according to the relationship $c = \sqrt{E / \rho}$.

It is noted that the speed of sound increases with the stiffness of the material, and decreases with the density. In solids, the velocity of sound depends on density of the material, not its temperature. Solid materials, such as steel, conduct sound much faster than air.

Propagation speeds for other media are given in figure 22. For each degree Centigrade increase in temperature, the speed of sound increases by 0.61 m/sec. It appears apparent from the data available, that the precision achieved by field determinations, laboratory studies, and calculation of the velocity of sound is adequate.

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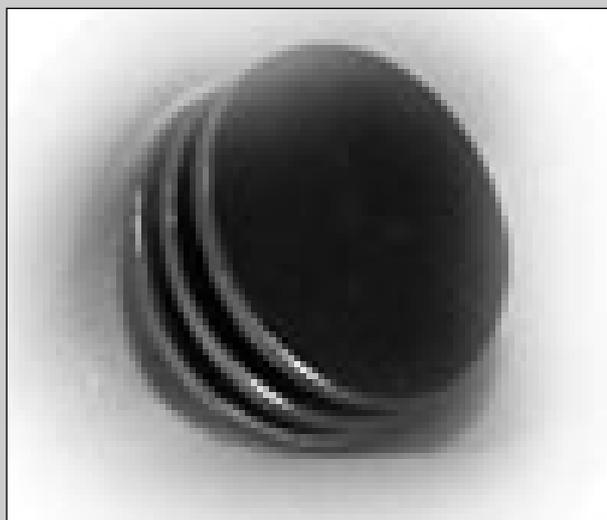
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